Modelling of Radio-frequency Heating of Piles of *Pinus radiata* Wood

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The present work studied the numerical modelling of heat transfer in a pile of *Pinus radiata* samples of a square cross-section by using radio frequency heating. More precisely, the study focused on the effects of the energy transferred to a dielectric material (wood) from an electromagnetic field, which required the calculation of the dielectric loss factor and its correlation with conservative equations. In this way, the temperature distribution across *Pinus radiata* samples was obtained through the integration of the energy equation using the finite volume method. The numerical results were compared to experimental data obtained from three experiments of radio-frequency heating of wood samples of a crosssectional area of 4 in \times 4 in, 3.1-m-long, and 20 cm of separation between plates. According to the observed linear behavior of the heat transfer process, the numerical results of the transient variation of temperature were in agreement with the experimental data.

Keywords: Dielectric heating; Dielectric loss factor; Heat treatment; Pinus radiata; Temperature simulation

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INTRODUCTION

Pinus radiata D. Don is one of the most important wood species in the Chilean wood industry (Infor 2016). *P. radiata* is often used as packing material, which requires it to be sterilized through vapour to fulfill phytosanitary requirements (Ananías *et al.* 2013). One alternative to using vapour to sterilize piles of *Pinus radiata* wood is to heat it through radio frequency electromagnetic radiation (Esquivel-Reyes *et al.* 2017). In contrast, dielectric heating, which implies the implementation of technologies such as radio frequency and microwaves, is used in several industries related to sterilization and drying of inert and organic materials (Wang and Wang 2008).

In traditional heating and drying processes, the thermal resistance of wood logs and woodpiles limits heat and mass transfer. Therefore, the slowest heating or drying regions are found in the thermal center of the treated product, which leads to over-processing of surface regions. However, radio-frequency heating is dependent upon the product itself, its geometry, and its thermophysical properties. As radio-frequency electromagnetic waves heat volumetrically, they can theoretically improve heating uniformity (Misra *et al.* 2015).

The heating process of wood for sterilization implies the loss of moisture to a lesser degree. Nevertheless, this phenomena could be important, depending on the performed study (Huang *et al.* 2013). To perform the modelling of the sterilization process of wood, both the heat and mass transfers must be considered (Jia *et al.* 2017).

Radio-frequency radiation consists of electromagnetic waves and represents a nonionizing type of radiation that causes molecular motion through migration of ions and rotation of dipoles without altering the molecular structure. A material can be heated by applying energy to it in the form of electromagnetic waves (Orfeuil 1987). The heating effect originates from the ability of an electric field to exert a force on charged particles. The heating effect depends on the frequency and the applied power of the electromagnetic waves, as well as on the material properties.

The dielectric properties of materials are defined by two parameters, the dielectric constant and the dielectric loss factor. The dielectric constant describes the ability of a molecule to be polarized by an electric field. The dielectric loss factor measures the efficiency with which the energy of the electromagnetic radiation is converted into heat. The ratio of the dielectric loss factor to its dielectric constant is expressed as the dissipation factor (Haswell and Kingston 1997). While electromagnetic waves penetrate the sample, energy is absorbed depending on the dissipation factor of the material. The greater the dissipation factor of a sample, the less electromagnetic energy will penetrate into it because it is rapidly absorbed and dissipated. Radio-frequency electromagnetic energy is transformed into heat via two mechanisms that occur simultaneously: ionic conduction and dipole rotation. The dissipation factor changes with temperature, which affects ion mobility and concentration. The second mechanism (dipole rotation) refers to the alignment of molecules that have a permanent or induced dipole moment. The relative contribution of each energy conversion mechanism depends mainly on the temperature. For small molecules, such as water, the contribution of a dipole rotation decreases as the sample temperature increases, whereas the contribution of ionic conduction increases as the product temperature is increased. Energy delivery to the biomaterial by radio-frequency radiation depends on a number of variables such as power input, exposure time, material properties, and the size of the sample. However, the greater the dissipation factor of a sample, the less it will be penetrated by the electromagnetic wave. Hence, in large samples with high dissipation factors, heating of the central part may still depend on thermal conductance (Misra et al. 2015).

From the modelling point of view, two approaches are found. The first one is focused on the effects of the electromagnetic field on the heat transfer to wood, whose dielectric parameters are known, particularly the dielectric loss factor, which are correlated through conservation equations (Koumoutsakos *et al.* 2001; Jia *et al.* 2015). The second approach is based on the description of the electric and magnetic fields according to the Maxwell equations (Bucki and Perré 2003, Marra *et al.* 2009). The present study used the former approach. Therefore, the dielectric loss factor was understood as the wood capacity as dielectric material of transforming electrical energy into thermal energy (Resch 2006; Avramidis 2016).

Particularly, a methodology for determining the dielectric loss factor from heating curves is presented, which is useful to determine the required power of the energy source in a differential mathematical model. The main goal of the present work is the numerical calculation of the spatial and transient distribution of temperature during the sterilization of moist *P. radiata* D. Don wood by using radio-frequency heating.

EXPERIMENTAL

Materials

The radio-frequency (RF) heating process was performed in triplicate on wood piles of square section samples measuring 4 in \times 4 in and 3.1-m-long (total volume of 3 m³). The *Pinus radiata* D. Don wood used was 14-years-old and obtained from the Bío Bío region in Chile. The samples of green wood were collected from the inner wood of randomly selected logs.

The separation between plates was kept at 200 mm. The wood load was mounted in the RF equipment (Saga HF-VD30SA, Shijiazhuang, Hebei, China), as shown in Fig. 1. The work power and the oscillation frequency were set at 30 kW and 6.78 MHz, respectively.



Fig. 1. Radio-frequency equipment: 1) High frequency (HF) generator, 2) Plate electrodes, 3) Autoclave, 4) Hydraulic system, 5) Vacuum pump, 6) Cooling tank, 7) Water pump, 8) Cooling tower, 9) Condensate tank, 10) Load cells, and 11) Trolley (Torres 2017)

Temperatures during heating were collected every 10 s using a data acquisition equipment and optical fiber temperature sensors (Oriental Rayzer, ORZ-FTM1000, Beijing, China) distributed as is shown in Fig. 2.



Fig. 2. Wood load (4 in \times 4 in \times 3.1 m): a) position of temperature sensors between plates (20 cm), and b) detail of the position of the sensors inside the wood load

Methods

Phenomenological modelling

The first implemented mathematical model was found in the adsorbed thermal energy due to the presence of an electromagnetic field, shown in Eq. 1,

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$$\Delta E_{System} = q_{\varepsilon} \tag{1}$$

which describes an energy balance between the heat flow generated by RF (q_{ε}) and the variation of the energy of the system (ΔE_{System}). If a transformation between electromagnetic energy and sensible thermal energy is considered (dielectric heating) without heat loss, and the inner conductive thermal resistance is considered to be low, the following ordinary differential equation that describes the transient variation of the average temperature of wood is obtained (Wang *et al.* 2005) in Eq. 2,

$$\frac{d(\rho c_p T)}{dt} = c_0 f E^2 \varepsilon'' \tag{2}$$

where c_p is the specific heat (J/kg°C), ρ is the apparent density of wood (kg/m³), *E* is the intensity of the electric field (V/m), *f* is the frequency (Hz), ε " is the dielectric loss factor, and c_0 is a constant equal to 55.33 × 10⁻¹². In this study, saturated wood with a homogeneous humidity content was considered, and the loss of humidity was not taken into account, therefore the differential equation may be written as,

$$\frac{dT}{dt} = c_0 \frac{f E^2 \varepsilon''}{\rho c_p} = c_1 \tag{3}$$

which may be solved analytically using Eqs. 4 and 5,

$$T(t) = T_0 + c_1 t \tag{4}$$

$$c_1 = c_0 \frac{f E^2 \varepsilon''}{\rho c_p} \tag{5}$$

which represents a linear increasing of temperature as a function of time.

RESULTS AND DISCUSSION

The experimental results for the transient variation of temperature obtained during the dielectric heating of wood are shown in Fig. 2.





The results shown in Fig. 2 clearly demonstrate that there was a linear increase in the average temperature of the wood in all tests. Moreover, it is important to note that the average temperature of all tests (T_pro_4x4) had a linear coefficient of determination almost equal to one ($R^2 = 0.9998$). These experimental results validate the mathematical model represented by Eq. 2, which solution is given by Eq. 4. This solution is valid when there is no heat losses and the inner thermal resistance is negligible. Consequently, and according to Zhou and Avramidis (1999), the average effective power transferred to the wood load is given by Eq. 6,

$$P_{transferred} = V \times \rho \times c_p \times c_1 \tag{6}$$

Then, considering the values for density and specific heat of *P. radiata* of 800 kg/m³ and 2.2 J/kg°C, respectively (Acuña 2017), and the c_1 value calculated from Eq. 5, the following result is obtained in Eq. 7:

$$P_{transferred} = 3 \times 8002.2 \times 0.0032 = 16.9 \, kW \tag{7}$$

Finally, if the transferred power and the dielectric field parameters (f and E) are known, the dielectric loss factor can be obtained through Eq. 8:

$$\varepsilon'' = \frac{c_o P_{transferred}}{fE^2} \tag{8}$$

One-dimensional Model

A second approach is based on the description of the spatial distribution of temperature. Therefore, mathematical models based on differential transport equations that describe the variation of the dependent variable in a one-dimensional (1D), two-dimensional (2D), or three-dimensional (3D) space is required. The available temperature data were collected along the horizontal symmetry axis, located in the cross-sectional plane in three characteristic points: at the surface, middle, and center of the wood load. Therefore, a one-dimensional heating model was considered using the previously obtained values of the transferred power ($P_{transferred}$) and the dielectric loss factor (ε "), when the inner thermal resistance was not considered. Under these considerations, the thermal energy balance of wood is given by Eq. 9,

$$\Delta E_{System} = \nabla(q_{con}) + q_{\varepsilon}(x) \tag{9}$$

where the equation represents an energy balance that considers the conductive heat transfer flow q_{con} (W/m²), a generation term $q_{\varepsilon}(x)$ in function of the independent variable (x), which corresponds to the gained heat due to the electric field, both terms induce variation of the energy rate of the system ΔE_{System} . The generation term, a product of the transformation of electromagnetic energy in sensible thermal energy (dielectric heating), was added to the conductive heat transfer according to the Fourier's Law, therefore Eq. 10 may be written as,

$$\rho c_{\rho} \frac{\partial T}{\partial t} = \frac{\partial}{\partial x} \left(k \frac{\partial T}{\partial x} \right) + q_{\varepsilon}(x) \tag{10}$$

where *T* is the temperature (°C), *t* is the time (s), *x* is the distance along the thickness of the wood pile (m), ρ is the density (kg/m³), c_p is the specific heat (J/kg °C), *k* is the thermal conductivity (W/m °C), and $q_{\varepsilon}(x)$ is the transferred heat from the electromagnetic field (W/m³). Equation 11 is under the following initial and boundary conditions,

$$T = T_0, \text{ for } t = 0$$

$$\pm k \frac{dT}{dx} = h(T_{\infty} - T_a) \text{ at } x = \pm a$$
(11)

where T_0 is the initial temperture (°C), T_{∞} is the environment temperature (°C), T_a is the surface temperature of Wood (°C), and *h* is the convective coefficient (W/m² °C).

Before solving Eq. 10 according to the established conditions in Eq. 11, the values of h and $q_{\varepsilon}(x)$ are required. The determination of h may be performed with the consideration that the energy from the wood to the inner side of the chamber of the radiofrequency equipment is equal to the transferred thermal energy from the inner side of the chamber to the environment (Q_{env}). Therefore,

$$h_{wood} = \frac{Q_{env}}{A_{wood}(T_{\infty} - T_{a})} \tag{12}$$

where A_{wood} is the surface area of wood (m²), T_{∞} as the ambient temperature (°C), and,

$$Q_{env} = \frac{A_{Chamber}\Delta T_{env}}{R_{eq}} \tag{13}$$

where ΔT_{env} is the variation of temperature (°C) between the inner and outer side of the chamber, A is the heat transfer area (m²), and R_{eq} is the equivalent thermal resistance (°C/W), given by Eq. 14,

$$R_{eq} = \left[\frac{1}{A_i h_i} + \frac{L}{A k_{wall}} + \frac{1}{A_e h_e}\right] \tag{14}$$

where k_{wall} is the thermal conductivity of the wall of the chamber of the radio-frequency equipment (W/m °C), *L* is the thickness of the chamber (m), and subscripts (i) and (e) refer to the inner and outer side of the chamber, respectively. The term $q_{\varepsilon}(x)$ is calculated according to the described preceding procedure for the temperature values close to the surface, the center, and the middle point between the center and the surface. From the first test (Fig. 3b), the obtained values of the temperature are shown in Fig. 3a.



а



Fig. 3. Transient variation of temperature of the wood (a) and experimental values for the thermal energy source term according to the position (b).

After the numerical integration of Eq. 10 using the finite volume method, which is described by Salinas *et al.* 2015, the temperature distribution shown in Fig. 4 was obtained.



Fig. 4. Spatial variation of temperatures obtained experimentally (discrete symbols) and numerically (segmented line): $\Delta x = 0.025$ m, and $\Delta t = 10$ s

The results showed a perfect correlation between the experimental and simulated data for $\Delta x = 0.025$ (m), and $\Delta t = 10$ (s). These numerical parameters were determined according to convergence analysis for numerical algorithms to solve the partial differential equations.

CONCLUSIONS

- 1. The measured spatial and transient variation of the average temperatures of a wood load during radio-frequency heating showed a linear behavior and therefore may be easily simulated by a linear mathematical model.
- 2. The variation of the temperature between the surface and the center followed the non-uniformity of the electric field and to the anisotropy of the dielectric material (wood).
- 3. The heating curves may be used for the determination of the dielectric properties, specifically the dielectric loss factor.
- 4. The spatial variation of the temperature may be modelled through the diffusive transport heat equation, considering a heat source that describes the transferred thermal energy from the electromagnetic field.
- 5. Given the linear behavior of the radio-frequency heating process, the simulated values of the transient variation of temperature were in agreement with the experimental values.

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